

# The Area Vortex for Modeling Flow through Smoothly Heterogeneous Aquifers

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## Motivation

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- Spatial heterogeneity is one of the major remaining hurdles in the development of AEM
- Tiling the domain with inhomogeneities is computationally expensive
- Methods can be developed to model non-piecewise constant heterogeneity without tiling

## Outline

- A conversation about modeling continuously variable conductivity
- The multi-quadric area vortex
  - Incremental advance
  - Not particularly useful
- Potential-based formulations
  - Some crazy ideas
  - "stir the pot"

## The Multi-quadric Area Vortex: Governing Equations

Water balance (in terms of stream function):

$$\frac{\partial}{\partial x} \left( \frac{1}{k} \frac{\partial \Psi}{\partial x} \right) + \frac{\partial}{\partial y} \left( \frac{1}{k} \frac{\partial \Psi}{\partial y} \right) = 0$$

Re-expressed as the Laplacian of the stream function:

$$\frac{\partial^2 \Psi}{\partial x^2} + \frac{\partial^2 \Psi}{\partial y^2} = -N_\Psi = \frac{1}{k} \frac{\partial k}{\partial x} \frac{\partial \Psi}{\partial x} + \frac{1}{k} \frac{\partial k}{\partial y} \frac{\partial \Psi}{\partial y}$$

$$\nabla^2 \Psi = \left[ \frac{\partial \ln k}{\partial y} Q_x - \frac{\partial \ln k}{\partial x} Q_y \right]$$

Specified

Dependent upon all elements

## The Multi-quadric Area Vortex

- Similar in concept to the multi-quadric area sink (Strack & Jankovic 1999)
  - Operates on the stream function rather than discharge potential
  - Laplacian of the stream function (related to curl), rather than Laplacian of the potential function (leakage), is interpolated using multi-quadric radial basis functions

$$N_\psi = \sum a_n r_n$$

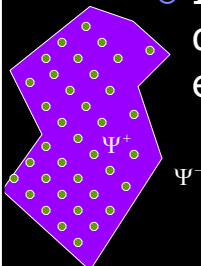
- Stream function Laplacian,  $N_\psi$ , calculated from flow solution and gradients in conductivity at control points:

$$N_\psi = \frac{1}{k} \frac{\partial k}{\partial y} Q_x - \frac{1}{k} \frac{\partial k}{\partial x} Q_y$$

- Requires modified BCs along border

## The Multi-quadric Area Vortex: Internal functional form

- Multi-quadric coefficients ( $a_n$ ) chosen to meet curl in a least squares sense at a set of control points
- Internal stream function calculated directly from MQ coefficients, external function=0



$$\nabla^2 \Psi^+ = -N_\psi$$

$$\nabla^2 \Psi^- = 0$$

# The Multi-quadric Area Vortex: Boundary Conditions

- Uses 3 line element types along border

- Line doublet:

- Removes jump in stream function  $\Psi^+ = \Psi^-$

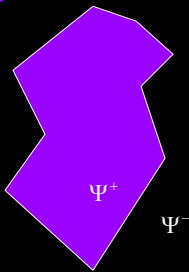
- Line vortex:

- Removes jump in tangential component of flow due to interior stream function
    - Subject to net curl of zero around vortex border

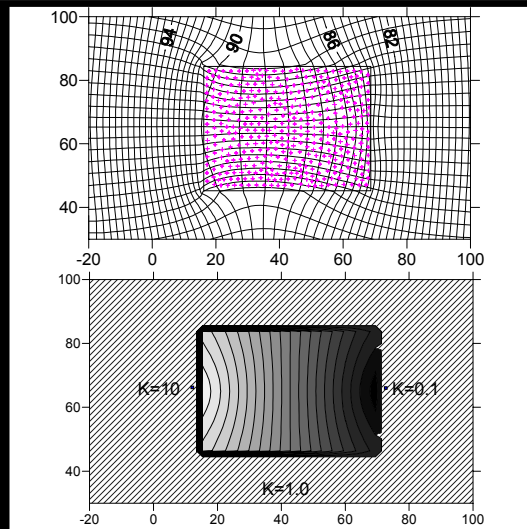
- Line dipole:

- Maintains law of refraction along element border (specific to area vortex for variable  $k$ )

$$\frac{Q_t^+}{k^+} = \frac{Q_t^-}{k^-} \quad Q_n^+ = Q_n^-$$



# The Multi-quadric Area Vortex: Simple test case



## The Multi-quadric Area Vortex: Difficulties



- Since potential is undefined within area vortex, numerical integration is required to calculate head
  - Internal Dirichlet boundaries problematic
- *Many* RBFs required to accurately resolve stream function divergence
  - Singularities in flow or conductivity within AV worsen issues
  - Prohibitive requirement for complex systems
- Can be highly non-linear
  - Iterative solution required
  - Computationally expensive for mildly complex flow systems

## Alternative Formulation



- Use an artificial potential w.r.t. reference  $k$ :

$$\Phi^* = \begin{cases} \frac{1}{2} \bar{k} \phi^2 & \text{Potential defined! -} \\ \bar{k} H^2 - \frac{1}{2} \bar{k} H \phi & \text{head available!} \end{cases}$$

- Discharge no longer the derivative of this "potential" function ( $Q_x = -k' d\Phi^*/dx$ ;  $k = k' \bar{k}$ )

- Governing Eqn. in terms of revised potential:

$$\nabla^2 \Phi^* = - \frac{\partial Y'}{\partial x} \frac{\partial \Phi^*}{\partial x} - \frac{\partial Y'}{\partial y} \frac{\partial \Phi^*}{\partial y}$$

- $Y'$  is the perturbation in log conductivity

$$Y' = \ln\left(\frac{k}{\bar{k}}\right) = \ln(k')$$

## Alternative Formulation: Radial Coordinate System

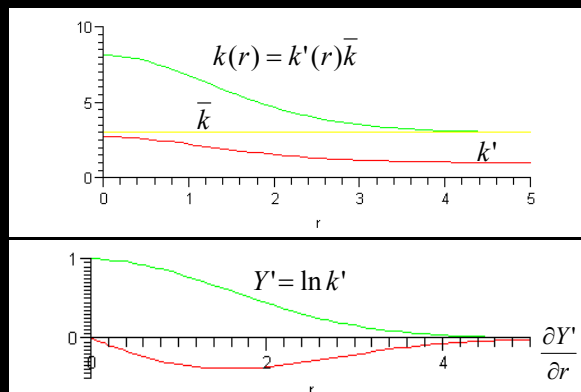
- Governing Equation:

$$\nabla^2 \Phi^* = - \frac{\partial Y'}{\partial r} \frac{\partial \Phi^*}{\partial r}$$

- $Y'(x,y)$  (and therefore  $k(x,y)$ ) represented by superposition of radial basis functions

- $Y' = a * F(r)$  where  $F(0) = 1$ ;  $F(\infty) = 0$ 
  - $Y' = a * \exp(-br^n)$  (exponential/Gaussian)
  - $Y' = a * b / (b + r^n)$  (Monod/Langmuir)

## Alternative Formulation : Radial Coordinate System



## Alternative Formulation: Elementary Solutions



- Solution may be superimposed upon standard Harmonic potential solution

$$\nabla^2\Phi^* + \nabla^2\Phi = -\frac{\partial Y'}{\partial r} \frac{\partial(\Phi^* + \Phi)}{\partial r} + 0$$

- Each harmonic element supplemented with additional curl term for each radial basis function
  - Standard solutions must be derived for basic analytic elements
    - Difficulty differs for each RBF Functional form
    - Wells and uniform flow are feasible
    - Line elements may not be amenable to closed-form solution

## Alternative Formulation: Advantages



- Potential, and therefore head, still defined
  - Consistent with existing 2D AEM solutions when  $dk/dx = dk/dy = 0$ ;
  - No numerical integration necessary
  - Governing eqn. can be solved with existing MQ area sinks
- Exact- no interpolation necessary, though functional approximations may be required
- Easily extended to 3D
- Potentially fast- added degrees of freedom (DOFs) for the most part directly dependent upon existing DOFs

## Alternative Formulation: Disadvantages



- Mathematically difficult
  - Initial derivations indicate that obtaining even simple solutions will be challenging
- Particular solutions are singular for useful RBFs, and will have to cancel out
- Requires  $M^*(N+M+1)$  degrees of freedom for  $N$  harmonic element coefficients and  $M$  radial basis functions (memory++)

## $k(x,y)$ : Other approaches



- Georgitza, 1969; Halek & Svec, 1979
- Define a function  $\Phi^* = \sqrt{k'} \cdot \bar{k} H \phi$      $k(x,y) = k'(x,y) \bar{k}$
- Discharge obtained from  $Q_x = -\sqrt{k'} \frac{d\Phi^*}{dx} - \Phi^* \frac{d\sqrt{k'}}{dx}$

General  
Inhomogeneous  
Medium

$$k'(x,y) = F(x,y)$$

$$\nabla^2 \Phi^* = \frac{\Phi^*}{\sqrt{k'}} \cdot \nabla^2 \sqrt{k'}$$

Helmholtz  
Inhomogeneous  
Medium

$$\nabla^2 \sqrt{k'} = \alpha^2 \sqrt{k'}$$

$$\nabla^2 \Phi^* = \alpha^2 \Phi^*$$

Harmonic  
Inhomogeneous  
Medium

$$\nabla^2 \sqrt{k'} = 0$$

$$\nabla^2 \Phi^* = 0$$

Piecewise  
Homogeneous  
Medium

$$\sqrt{k'} = 1$$

$$\nabla^2 \Phi^* = \nabla^2 \Phi = 0$$



## Conclusions

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- There is still much to be done to effectively simulate continuously varying properties with the analytic element method
- The multi-quadric area vortex, while conceptually similar to the MQ area sink, suffers from significant drawbacks
- Alternative approaches in terms of artificial potential seem much more promising